

The Blow-Up Time For Reaction-Diffusion Equations With Dirichlet Boundary Conditions

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Abstract - This paper is concerned with the study of the numerical approximation of the following initial-boundary value problem:

$$\left\{ \begin{array}{l} u_t - u_{xx} - \frac{a}{x} u_x = f(u(x, t)), 0 < x < 1, t \in (0, T), \\ u(0, t) = 0, t \in (0, T) \\ u(1, t) = 0, t \in (0, T) \\ u(x, 0) = u_0(x), 0 \leq x \leq 1 \end{array} \right.$$

Where $a \geq 0$, and $f(u) = u^q$, where $q > 0$.

We obtain some conditions under which the positive solution of a semi-discrete form of the above problem blows up in a finite time and estimate its semi-discrete blow up time. Under some assumptions, we also show that the semi-discrete blow up time converges to the real one when the mesh parameter goes to zero. Finally, we give some numerical results to illustrate our analysis.

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I. INTRODUCTION

Consider the following initial boundary value problem:

$$u_t - u_{xx} - \frac{a}{x} u_x = f(u(x, t)), 0 < x < 1, t \in (0; T) \quad (1)$$

$$u(0, t) = 0, t \in (0, T) \quad (2)$$

$$u(1, t) = 0, t \in (0, T) \quad (3)$$

$$u(x, 0) = u_0(x), 0 \leq x \leq 1 \quad (4)$$

Which models the temperature distribution of a large number of physical phenomenon from physics, chemistry and biology. It is known that a lot of physical and biological phenomena can be described by the following reaction-diffusion equation:

$$u_t(t, x) = u_{xx}(t, x) + f(u(x, t)) \quad Eq \ 1$$

Where the variable $u(t, x)$ can be seen as the temperature in a chemistry reaction or the population density of a biological species (see [21]).

The nonlinearity of f stands for a net heat source (the natural heat source subtract the dissipative heat source), which satisfies some growth conditions: it becomes positive when u is large enough, or is positive for all $u > 0$.

The second order derivative u_{xx} represents the diffusion. In a chemical reaction, we know that the chemical reaction will generate heat and the high temperature will accelerate the chemical reaction, thus the temperature will become higher and higher. Therefore, if the initial temperature is high enough, then the temperature will likely become high in a finite time, which is called a blow-up phenomenon.

Similarly, in a biological species, the large population density will bring the high birth rate, on the other hand; the high birth rate will accelerate the population growth. Hence, if the initial population density is large enough, then blow-up will happen too.

Motivated by such phenomena arising in various applied fields, the blow-up problem for Eq1 has been studied by many authors (cf. [13], [22] and references therein). For example, the classical papers ([15], [17]) consider the Cauchy problem:

$$u_t = u_{xx} + au + u^p (x \in R, t > 0), u(0, x) = u_0(x) \geq 0 (x \in R) \quad Eq2$$

, with $a = 0$ and prove that, when $1 < p \leq 3$, the nontrivial solution must blow-up in a finite time; when

$p > 3$, this problem has time-global solutions for some small initial data. In [2], [8], [23], the authors studied this problem with $a < 0$ and gave trichotomy result: any solution either blows up in a finite time, or vanishes (i.e., $u \rightarrow 0$ as $t \rightarrow \infty$), or converges to an evenly decreasing positive steady state as $t \rightarrow \infty$.

The author in [18] studied the corresponding steady stats problem on annulus with $a < 0$. Moreover in [26], the author studied the corresponding steady stats problem of Eq2, with $a \geq 0$ and establishes some existence and nonexistence results. The theoretical study of blow-up solutions for semi linear parabolic equations has been the subject of investigations of many authors (see [3],[6],[13],[16],[20],[27] and the references cited therein). The authors have proved that under some assumptions, the solution of (1)-(4) blows up in a finite time and the blow-up time is estimated.

Let us notice that if we consider the semi linear equation

$$u_t = u_{xx} + u^p, x \in B, t \in (0, T),$$

With boundary conditions

$$u(x, t) = 0, x \in S, t \in (0, T),$$

And initial data

$$u(x, 0) = u_0(x) \in \bar{B},$$

Where $B = \{x \in R^n : |x| < 1\}$,

$S = \{x \in R^n : |x| = 1\}$, the radial symmetric solutions are solutions of (1)-(4) with $a = 0$ and $f(u(x, t)) = u^p, p = n - 1$.

In this paper, we are interested in the numerical approximations of the problem (1) – (4) using a semi-discrete form of (1) – (4). Here $(0, T)$ is the maximal time interval on which the solution exists. The time T may be finite or infinite. When T is infinite, we say that the solution u exists globally. When T is finite, the solution u develops a singularity in a finite time, namely :

$$\lim_{t \rightarrow T} \|u(x, t)\|_{\infty} = +\infty.$$

Where $\|u(x, t)\|_{\infty} = \max_{0 \leq x \leq 1} |u(x, t)|$. In this case, we say that the solution u blows up in a finite time and the time T is called the blow-up time of the solution u .

II. SEMIDISCRETE PROBLEM

Let I be a positive integer and define the grid

$$x_i = ih, 0 \leq i \leq I, \text{ where } h = 1/I. \text{ Approximate}$$

the solution u of (1) – (4) by the solution

$$U_h(t) = (U_0(t), U_1(t), \dots, U_I(t))^T$$

$$+ \delta U_i(t) + U_i^q(t), 1 \leq i \leq I - 1, t \in (0, T_b^h)$$

$$\frac{dU_i(t)}{dt} = \delta^2 U_i(t) + \frac{a}{ih} \delta^i, \quad (5)$$

$$U_0(t) = 0, t \in (0, T_b^h), \quad (6)$$

$$U_I(t) = 0, t \in (0, T_b^h), \quad (7)$$

$$U_i(0) = \varphi_i, 0 \leq i \leq I, \quad (8)$$

Where

$$\delta^2 U_i(t) = \frac{U_{i+1}(t) - 2U_i(t) + U_{i-1}(t)}{h^2}, 1 \leq i \leq I - 1$$

$$+ \delta U_i(t) = \frac{U_{i+1}(t) - U_i(t)}{h}, 1 \leq i \leq I - 1.$$

Here, $(0, T_b^h)$ is the maximal time interval on which

$\|U_h(t)\|_{\infty}$, is finite, where

$$\|U_h(t)\|_{\infty} = \max_{0 \leq i \leq I} |U_i(t)|.$$

When T_b^h is finite, we say that the solution

$U_h(t)$ blows-up in a finite time and the time

T_b^h is called the blow-up time of the solution

$U_h(t)$. We give some conditions under which the

solution of (5)-(8) blows up in a finite time and estimate its semidiscrete blow-up time. We also show that the semidiscrete blow-up time converges to the theoretical one when the mesh size goes to zero. A similar study has been undertaken in [1],[10],[12] and [14]. In [1] and [14], the authors have considered the equation (1) for $a = 0$ and $f(u(x, t)) = u^q$ with Dirichlet boundary conditions and nonnegative initial data. Our work was also motivated by the papers in [1], [12] and [14]. In [1], the authors have considered the problem (1)–(4) in the case where the parameter $a = 0$, and the source is u^q . They have proved that the solution of the semi discrete scheme (5)-(8) blows up in a finite time and its semi discrete blow-up time converges to real one when the mesh size goes to zero in the case where the initial data is symmetric and large enough.

In [19], the author has shown that the solution of a discrete form of the equation $u_t = u_{xx} + u^2$ with Dirichlet boundary condition and large initial data blows up in a finite time and the discrete blow-up time converges to the real one when the mesh parameter goes to zero. In [3], semi discrete and discrete schemes have been used to study the phenomenon of extinction (we say that a solution extincts in a finite time if it reaches the values zero in a finite time). Numerical methods for heat equations with nonlinear boundary conditions have been described in [12], and [14].

The paper is organized as follows. In the next section, we give some properties concerning our scheme. In the fourth section, under some conditions, we prove that the solution of the semidiscrete problem (5)-(8) blows up in a finite time and estimate its semidiscrete blow-up time. In the fifth section, we study the convergence of the semidiscrete blow-up time. Finally, in the last section we report on some numerical experiments using several discretisations of (1)-(4).

III . PROPRIETE OF THE SEMIDISCRETE SCHEME

In this section, we give some results about the discrete maximum principle. The following lemma is a discrete form of the maximum principle.

Lemma 3.1

Let $c_h(t) \in C^0([0, T], R^{I+1})$ and $V_h(t) \in C^1([0, T], R^{I+1})$ such that for $t \in (0, T)$

$$+i V_i(t) + c_i(t) V_i(t) \geq 0, 1 \leq i \leq I-1$$

$$\frac{dV_i(t)}{dt} - \delta^2 V_i(t) - \frac{a}{ih} \delta^i, \quad t \in (0, T_b^h), \quad (9)$$

$$V_0(t) \geq 0, V_I(t) \geq 0 \quad (10)$$

$$V_i(0) \geq 0, 0 \leq i \leq I \quad (11)$$

Then we have $V_i(t) \geq 0, 0 \leq i \leq I, t \in (0, T)$.

Proof:

Let $T_0 < T$ and let $m = V_i(t)$. Since for $i \in \{0, \dots, I\}$, $V_i(t)$ is a continuous function, there exists $t_0 \in [0, T_0]$ such that $m = V_{i_0}(t_0)$ for a certain $i_0 \in \{0, \dots, I\}$. If $i_0 = 0$ or $i_0 = I$, we have $m > 0$. When i_0 is between 1 and $I-1$, we observe that :

$$\frac{dV_{i_0}(t_0)}{dt} = \lim_{k \rightarrow 0} \frac{V_{i_0}(t_0) - V_{i_0}(t_0 - k)}{k} \leq 0 \quad (12)$$

$$\delta^2 V_{i_0}(t_0) = \frac{V_{i_0+1}(t_0) - 2V_{i_0}(t_0) + V_{i_0-1}(t_0)}{h^2} \geq 0,$$

$$\text{if } 1 \leq i_0 \leq I-1 \quad (13)$$

$$+i V_{i_0}(t_0) = \frac{V_{i_0+1}(t_0) - V_{i_0}(t_0)}{h} \geq 0, \text{ if } 1 \leq i_0 \leq I-1 \quad (14)$$

Define the vector $Z_h(t) = e^{\lambda t} V_h(t)$ where λ is such that $c_{i_0}(t_0) - \lambda > 0$.

A straightforward computation reveals that

$$(i i_{i_0}(t_0) - \lambda) Z_{i_0}(t_0) \geq 0$$

$$+i Z_{i_0}(t_0) + i$$

$$\frac{dZ_{i_0}(t_0)}{dt} - \delta^2 Z_{i_0}(t_0) - \frac{a}{i_0 h} \delta^i$$

$$\text{if } 1 \leq i_0 \leq I-1 \quad (15)$$

We observe from (5)-(8) that

$$\frac{dZ_{i_0}(t_0)}{dt} \leq 0, \delta^2 Z_{i_0}(t_0) \geq 0, \quad \text{and}$$

$$+i Z_{i_0}(t_0) \geq 0$$

Using (9), we arrive at $(i i_{i_0}(t_0) - \lambda) Z_{i_0}(t_0) \geq 0$, which implies that $Z_{i_0}(t_0) \geq 0$

Therefore $V_{i_0}(t_0) = m \geq 0$ and we have the desired result, and the proof is complete.

Another version of the discrete maximum principle is the following comparison lemma.

Lemma 3.2

Let $V_h(t), U_h(t) \in C^1([0, T], R^{I+1})$ and

$g \in C^0(R \times R, R)$ such that for $t \in (0, T)$

$$+i U_i(t) + g(U_i(t), t), 1 \leq i \leq I-1 \quad (16)$$

$$+i V_i(t) + g(V_i(t), t) < \frac{dU_i(t)}{dt} - \delta^2 U_i(t) - \frac{a}{ih} \delta^i$$

$$\frac{dV_i(t)}{dt} - \delta^2 V_i(t) - \frac{a}{ih} \delta^i$$

$$V_0(t) < U_0(t), V_I(t) < U_I(t), \quad (17)$$

$$V_i(0) < U_i(0), 0 \leq i \leq I \quad (18)$$

Then we have $V_i(t) < U_i(t), t \in (0, T)$.

Proof:

Define the vector $Z_h(t) = U_h(t) - V_h(t)$. Let t_0 be the first $t > 0$ such that $Z_h(t) > 0$ for $t \in (0, t_0)$, but

$Z_{i_0}(t_0) = 0$ for a certain $i_0 \in \{0, \dots, I\}$. If $i_0 = 0$ or $i_0 = I$, we have a contradiction because of (17).

If i_0 is between 0 and $I-1$, we observe that

$$\frac{dZ_{i_0}(t_0)}{dt} = \lim_{k \rightarrow 0} \frac{Z_{i_0}(t_0) - Z_{i_0}(t_0 - k)}{k} \leq 0$$

$$\delta^2 Z_{i_0}(t_0) = \frac{Z_{i_0+1}(t_0) - 2Z_{i_0}(t_0) + Z_{i_0-1}(t_0)}{h^2} \geq 0, \text{ if } 1 \leq i_0 \leq I-1$$

$$+i Z_{i_0}(t_0) = \frac{Z_{i_0+1}(t_0) - Z_{i_0}(t_0)}{h} \geq 0, \text{ if } 1 \leq i_0 \leq I-1$$

Which implies that

$$+\dot{Z}_{i_0}(t_0)+g(U_{i_0}(t_0),t_0)-g(V_{i_0}(t_0),t_0)\leq 0$$

$$\frac{dZ_{i_0}(t_0)}{dt}-\delta^2 Z_{i_0}(t_0)-\frac{a}{ih}\delta^i, \quad \text{if}$$

$$1\leq i\leq I-1$$

But this inequality contradict (16). This ends the proof.

IV. SEMIDISCRETE BLOWING-UP SOLUTIONS

In this section under some assumptions, we show that the solution U_h of the semidiscrete problem (5)-(8) blows up in a finite time and estimate its semidiscrete blow-up time. We need the following important lemma. This lemma at the first gives us a property of the operator δ^2 , then reveals a property of the operator $\dot{\delta}^i$.

Lemma 4.1

Let $U_h \in R^{I+1}$ such that $U_h \geq 0$. Then we have

$$\delta^2 f(U_i) \geq f'(U_i) \delta^2 U_i, \quad 0 \leq i \leq I$$

$$+\dot{Z}_{i_0} U_i$$

$$+\dot{Z}_{i_0} f(U_i) \geq f'(U_i) \delta^i, \quad 0 \leq i \leq I$$

Proof :

We apply Taylor's expansion to obtain :

$$f(U_{i+1}) = f(U_i) + (U_{i+1} - U_i) f'(U_i) + \frac{(U_{i+1} - U_i)^2}{2} f''(\theta)$$

, for $0 \leq i \leq I-1$

$$f(U_{i-1}) = f(U_i) + (U_{i-1} - U_i) f'(U_i) + \frac{(U_{i-1} - U_i)^2}{2} f''(\eta)$$

, for $1 \leq i \leq I$

where θ_i is an intermediate values between U_i and U_{i+1} , and η_i is an intermediate value between U_{i-1} and U_i .

These two equations above imply that

$$f(U_{i+1}) - 2f(U_i) + f(U_{i-1}) = (U_{i+1} - 2U_i + U_{i-1}) f'(U_i).$$

The first equation and the last one imply that :

$$\delta^2 f(U_i) = f'(U_i) \delta^2 U_i + \frac{(U_{i+1} - U_i)^2}{2h^2} f''(\theta_i) + \frac{(U_{i+1} - U_i)}{2h^2}$$

$$+\dot{Z}_{i_0} U_i + \frac{(U_{i+1} - U_i)^2}{2h} f''(\theta_i)$$

$$+\dot{Z}_{i_0} f(U_i) = f'(U_i) \delta^i$$

Using the convexity of f we obtain the desired result.

Lemma: 4.2

Let U_h be the solution of (5)-(8). Then we have

$$U_{i+1}(t) > U_i(t), \quad 0 \leq i \leq I-1, t \in (0, T_b^h), (19)$$

Proof.

Let t_0 be the first $t > 0$, such that $U_{i+1}(t) - U_i(t) > 0$ for $0 \leq i \leq I-1$, but

$$U_{i_0+1}(t) = U_{i_0}(t) \quad \text{for a certain } i_0 \in \{0, \dots, I-1\}$$

. Without lost of generality we may suppose that i_0 is the smallest integer which satisfies the above equality.

Now let us introduce the functions

$$Z_i(t) = U_{i+1}(t) - U_i(t) \quad \text{for } 0 \leq i \leq I-1, t > 0.$$

We get

$$\frac{dZ_{i_0}(t_0)}{dt} = \lim_{k \rightarrow 0} \frac{Z_{i_0}(t_0) - Z_{i_0}(t_0 - k)}{k} \leq 0$$

$$\delta^2 Z_{i_0}(t_0) = \frac{Z_{i_0+1}(t_0) - 2Z_{i_0}(t_0) + Z_{i_0-1}(t_0)}{h^2} \geq 0, \text{ if } 1 \leq i \leq I-1$$

$$+\dot{Z}_{i_0} Z_{i_0}(t_0) = \frac{Z_{i_0+1}(t_0) - Z_{i_0}(t_0)}{h} \geq 0, \text{ if } 0 \leq i \leq I-1$$

which implies that

$$+\dot{Z}_{i_0} Z_{i_0}(t_0) + f(U_{i_0+1}(t_0), t_0) - f(U_{i_0}(t_0), t_0) \leq 0$$

$$\frac{dZ_{i_0}(t_0)}{dt} - \delta^2 Z_{i_0}(t_0) - \frac{a}{i_0 h} \delta^i, \quad \text{if}$$

$$0 \leq i_0 \leq I-1$$

But this inequality contradicts (5)-(8) and we obtain the desired result.

Theorem 4.1

Suppose that there exists a positive integer A such that the initial data at (8) satisfies

$$+\dot{Z}_{i_0} U_i(0) + U_i^q(0) \geq A U_i^q(0)$$

$$\delta^2 U_i(0) + \frac{a}{ih} \delta^i, \quad 0 \leq i \leq I (20)$$

Then the solution $U_i(t)$ of (5)-(8) blows up in a finite time T_b^h and we have the following estimate

$$T_b^h \leq \frac{1}{A} \frac{\|U_h(0)\|_\infty^{1-q}}{(q-1)}$$

Proof

Since $(0, T_b^h)$ be the maximal time interval on which $\|U_h(t)\|_\infty$ is finite, our aim is to show that

T_b^h is finite and obeys the above inequality.

Introduce the following functions $f(U) = U^q$ and

J_h such that :

$$J_i = \frac{dU_i}{dt} - AU_i^q, 0 \leq i \leq I.$$

we shall prove that $J_h \geq 0, t \geq 0, 0 \leq i \leq I$.

By a straightforward computation we get:

$$\frac{dJ_i}{dt} - \delta^2 J_i = \frac{d}{dt} \left(\frac{dU_i}{dt} - \delta^2 U_i \right) - AqU_i^{q-1} \frac{dU_i}{dt} + A\delta^2 U_i^q, 0 \leq i \leq I.$$

From Lemma 4.1

$$\delta^2 f(U_i) \geq f'(U_i) \delta^2 U_i, \quad 0 \leq i \leq I$$

$$+ \zeta U_i$$

$$+ \zeta f(U_i) \geq f'(U_i) \delta^{\zeta} U_i, \quad 0 \leq i \leq I$$

we have $\delta^2 U_i^q \geq qU_i^{q-1} \delta^2 U_i$ which implies that

$$\frac{dJ_i}{dt} - \delta^2 J_i \geq \frac{d}{dt} \left(\frac{dU_i}{dt} - \delta^2 U_i \right) - AqU_i^{q-1} \left(\frac{dU_i}{dt} - \delta^2 U_i \right)$$

Use (5)-(8) to obtain

$$+ \zeta U_i$$

$$+ \zeta \frac{dU_i}{dt} - AqU_i^{q-1} \delta^{\zeta} U_i$$

$$\frac{dJ_i}{dt} - \delta^2 J_i \geq qU_i^{q-1} J_i + \frac{a}{ih} \zeta$$

$$0 \leq i \leq I. \quad (21)$$

Taking into account the expression of $J_h(t)$, we

$$\text{get } + \zeta \frac{dU_i}{dt} - \zeta, \text{ which implies that}$$

$$+ \zeta J_i = \delta^{\zeta} U_i$$

$$+ \zeta J_i + \zeta$$

$$+ \zeta \frac{dU_i}{dt} = \delta^{\zeta} U_i$$

$$+ \zeta J_i + AU_i^{q-1} \delta^{\zeta} U_i$$

Therefore, we get

$$\frac{dU_i}{dt} \geq \delta^{\zeta} U_i$$

$$+ \zeta J_i$$

$$+ \zeta U_i \geq \delta^{\zeta} U_i$$

band due to (20), we discover

$$\frac{dU_i}{dt} - AU_i^{q-1} \delta^{\zeta} U_i$$

that

$$+ \zeta \frac{dJ_i}{dt} \geq qU_i^{q-1} J_i, 0 \leq i \leq I$$

Obviously,

$$\frac{dJ_i}{dt} - \delta^2 J_i - \frac{a}{ih} \delta^{\zeta} U_i$$

$J_0(t) = 0, J_I(t) = 0$ and the hypotheses on the initial data ensure that $J_h(0) \geq 0$. It follows from Lemma 3.1 that $J_h(t) \geq 0$ for $t \in (0, T_b^h)$, which implies that

$$\frac{dU_i}{dt} \geq AU_i^q, \quad 0 \leq i \leq I.$$

This estimation may be rewritten as follows :

$$\frac{dU_i}{U_i^q} \geq Adt, \quad 0 \leq i \leq I \quad (22)$$

Integrating this inequality over $(0, T_b^h)$, we find that :

$$T_b^h - t \leq \frac{1}{A} \frac{U_i^{1-q}(t)}{(q-1)}, \quad 0 \leq i \leq I$$

(23)

which implies that $T_b^h \leq \frac{1}{A} \frac{\|U_h(0)\|_{\infty}^{1-q}}{(q-1)}$. Therefore

T_b^h is finite and the proof is complete.

Remark 4.1

The inequalities (22) imply that

$$T_b^h - t_0 \leq \frac{1}{A} \frac{\|U_h(t_0)\|_{\infty}^{1-q}}{(q-1)}, \quad \text{if } 0 \leq t_0 \leq T_b^h \text{ and}$$

there exists a constant $B > 0$ such that

$$U_i(t) \leq \frac{B}{(T_b^h - t)^{\frac{1}{q-1}}}, \quad 0 \leq i \leq I.$$

Theorem 4.2

Let $U_h(t)$ be the solution of (5)-(8). Then we have

$$T_b^h \geq \frac{\|U_h(0)\|_{\infty}^{1-q}}{(q-1)}.$$

Proof:

Let i_0 be such that $U_{i_0}(t) = \max_{0 \leq i \leq I} U_i(t)$. It is not hard to see that

$$\delta^2 U_{i_0}(t) = \frac{U_{i_0+1}(t_0) - 2U_{i_0}(t_0) + U_{i_0-1}(t_0)}{h^2} \leq 0, \text{ if } 1 \leq i \leq I-1$$

$$+ \zeta U_{i_0}(t) = \frac{U_{i_0+1}(t) - U_{i_0}(t)}{2h} \leq 0, \text{ if } 0 \leq i \leq I-1.$$

We deduce that $\frac{dU_{i_0}}{dt} \leq U_{i_0}^q$, which implies that

$\frac{dU_i}{U_i^q} \leq dt$. Integrating this inequality over $(0, T_b^h)$,

$$\text{we obtain } T_b^h \geq \frac{U_{i_0}^{1-q}(0)}{(q-1)}.$$

Use the fact that $U_{i_0}(0) = \|U_h(0)\|_\infty$, we obtain

$$T_b^h \geq \frac{\|U_h(0)\|_\infty^{1-q}}{(q-1)} \text{ and the theorem is completely proved.}$$

V. CONVERGENCE OF THE SEMIDISCRETE BLOW-UP TIME

In this section, we prove the convergence of the semidiscrete blow-up time. Under some assumptions, we show that the blow-up time of the solution for the semidiscrete problem converges to the real one when the mesh size tends to zero. In order to prove this result, firstly we show that for each fixed time interval $[0, T]$ where the solution u of (1)-(4) is defined, the solution $U_h(t)$ of (5)-(8) approximates u , when the mesh parameter h goes to zero by the following theorem.

Theorem 5.1

Assume that (1)-(4) has a solution $u \in C^{4,1}([0,1] \times [0, T])$ and the initial condition at (8) satisfies

$$\|\varphi_h - u_h(0)\|_\infty = o(1) \text{ as } h \rightarrow 0 \quad (24)$$

where $u_h(t) = (u(x_0, t), \dots, u(x_I, t))^T$. Then, for h sufficiently small, the problem (5)-(8) has a unique solution $U_h(t) \in C^1([0, T], R^{I+1})$ such that

$$\max_{0 \leq t \leq T} \|U_h(t) - u_h(t)\|_\infty = o(\|\varphi_h - u_h(0)\|_\infty + h) \text{ as } h \rightarrow 0 \quad (25)$$

Proof

Since $u \in C^{4,1}$, there exist positive constants K and M such that

$$\|u\|_\infty \leq K, q(K+1)^{q-1} \leq M \quad (26)$$

The problem (5)-(8) has for each h a unique solution $U_h(t) \in C^1([0, T_b^h], R^{I+1})$. Let $t(h)$ the greatest value of $t > 0$ such that

$$\|U_h(t) - u_h(t)\|_\infty < 1 \text{ for } t \in (0, t(h)). \quad (27)$$

The relation (24) implies that $t(h) > 0$ for h sufficiently small. Let $t^*(h) = \min(t(h), T)$. By the triangle inequality, we obtain

$$\|U_h(t)\|_\infty \leq \|u_h(t)\|_\infty + \|U_h(t) - u_h(t)\|_\infty, \text{ for } t \in (0, t^*(h)), \text{ which implies that}$$

$$\|U_h(t)\|_\infty \leq 1 + K \text{ for } t \in (0, t^*(h)). \quad (28)$$

Let $e_h = U_h(t) - u_h(t)$ be the error of discretization.

Since $u \in C^{4,1}$, using Taylor's expansion, we have for $t \in (0, t^*(h))$,

$$\begin{aligned} + \dot{e}_i(t) &= q \xi_i^{q-1} e_i + o(h) \\ \frac{de_i(t)}{dt} - \delta^2 e_i(t) - \frac{a}{ih} \delta^i, \end{aligned}$$

where ξ_i is an intermediate value between $U_i(t)$ and $u(x_i, t)$. Using (26) and (28), there exists a positive constant M such that

$$\begin{aligned} + \dot{e}_i(t) &\leq M |e_i(t)| + M h \\ \frac{de_i(t)}{dt} - \delta^2 e_i(t) - \frac{a}{ih} \delta^i, \\ 0 \leq i \leq I. \end{aligned} \quad (29)$$

Introduce the vector $Z_h(t)$ such that

$$Z_i(t) = e^{(M+1)t} (\|\varphi_h - u_h(0)\|_\infty + M h), \quad 0 \leq i \leq I.$$

A straightforward calculation reveals that

$$\begin{aligned} + \dot{Z}_i(t) &> M Z_i + M h \\ \frac{dZ_i(t)}{dt} - \delta^2 Z_i(t) - \frac{a}{ih} \delta^i, \quad 0 \leq i \leq I, \\ Z_I(t) &> e_I(t). \end{aligned}$$

It follows from Lemma 3.2 that $Z_h(t) > e_h(t)$ for $t \in (0, t^*(h))$.

In the same way, we also prove that

$Z_h(t) > -e_h(t)$ for $t \in (0, t^*(h))$, which implies that

$$\|U_h(t) - u_h(t)\|_\infty \leq e^{(M+1)t} (\|\varphi_h - u_h(0)\|_\infty + M h) \quad 0 \leq i \leq I.$$

Let us show that $t^*(h) = T$. Suppose that $T > t(h)$. From (27), we obtain

$$1 = \|U_h(t(h)) - u_h(t(h))\|_\infty \leq e^{(M+1)t(h)} (\|\varphi_h - u_h(0)\|_\infty + M h). \quad (30)$$

Since the term on the right hand side of the above inequality $1 \leq 0$ is impossible.

Consequently $t^*(h) = T$, and the proof is complete.

Now, we are in a position to prove the main theorem of this section.

Theorem 5.2

Suppose that the problem (1)-(4) has a solution u which blows up in a finite time T_b such that

$u \in C^{4,1}([0,1] \times [0, T])$ and the initial data at (4) satisfies $\|\varphi_h - u_h(0)\|_\infty = o(1)$ as $h \rightarrow 0$.

$V_h^{(n+1)} = (U_0^{(n+1)}, U_1^{(n+1)}, \dots, U_{I-1}^{(n+1)})^T$. It is not hard to see that $(A_h^{(n)})_{ii} \geq 0$, $(A_h^{(n)})_{ij} \leq 0, i \neq j$, $(A_h^{(n)})_{ii} > \sum_{i \neq j} |(A_h^{(n)})_{ij}|$.

These inequalities imply that the linear system has a unique solution for $n \neq 0$ and the discrete solution is nonnegative. For the proof, see for instance [3]. We need the following definition.

Definition 5.1

We say that the solution $U_h^{(n)}$ of (32)-(34) or (35)-(37) blows up in a finite time if $\lim_{n \rightarrow +\infty} |U_h^{(n)}| = +\infty$

and the series $\sum_{n=0}^{+\infty} \Delta t_n$ converges.

The quantity $\sum_{n=0}^{+\infty} \Delta t_n$ is called the numerical blow-up time of the solution $U_h^{(n)}$.

In the tables 1, 2, 3, 4, 5, 6, 7 and 8, we present the numerical blow-up times, values of n , the CPU times and the orders of the approximations corresponding to meshes of 16, 32, 64, 128. We take

for the numerical blow-up time $T^n = \sum_{j=0}^{n-1} \Delta t_j$ which is computed at the first time when $\Delta t_n = |T^{n+1} - T^n| \leq 10^{-16}$. The order(s) of the method is computed from

$$s = \frac{\log\left(\frac{T_{4h} - T_{2h}}{T_{2h} - T_h}\right)}{\log(2)}$$

Numerical experiments for

$$U_i^{(0)} = 20 \cos\left(\frac{i\pi h}{2}\right), q=2$$

First case: $a=0$

Table 1: Numerical blow-up times, numbers of iterations, CPU times (seconds), and orders of the approximations obtained with the explicit Euler method

l	t_n	n	CPU time	s
16	0.056343	7191	16	-
32	0.056231	27359	69	-
64	0.056210	103908	560	2.31
128	0.055203	405086	7052	1.68

Table 2: Numerical blow-up times, numbers of iterations, CPU times (seconds), and orders of the

approximations obtained with the implicit Euler method

l	t_n	n	CPU time	s
16	0.056376	7191	20	-
32	0.056240	27360	113	-
64	0.056213	103908	1460	2.22
128	0.055207	406009	20746	1.85

Second case: $a=1$

Table 3: Numerical blow-up times, numbers of iterations, CPU times (seconds), and orders of the approximations obtained with the explicit Euler method.

l	t_n	n	CPU time	s
16	0.067710	7406	13	-
32	0.066986	28264	70	-
64	0.066527	107691	594	0.76
128	0.066300	409301	20746	1.02

Table 4: Numerical blow-up times, numbers of iterations, CPU times (seconds), and orders of the approximations obtained with the implicit Euler method

l	t_n	n	CPU time	s
16	0.067802	7407	18	-
32	0.067008	28265	109	-
64	0.0665323	107692	1506	0.84
128	0.066303	409302	21000	1.05

Third case: $a=1.5$

Table 5: Numerical blow-up times, numbers of iterations, CPU times (seconds), and orders of the approximations obtained with the explicit Euler method

l	t_n	n	CPU time	s
16	0.076519	7555	13	-
32	0.075299	28888	69	-
64	0.074569	110314	590	0.85
128	0.074177	435971	8580	0.95

Table 6: Numerical blow-up times, numbers of iterations, CPU times (seconds), and orders of the approximations obtained with the implicit Euler method

l	t_n	n	CPU time	s
16	0.076700	7555	11	-
32	0.075332	28888	59	-
64	0.074577	110001	499	0.85
128	0.074177	420284	7580	0.95

Fourth case: $a=2$

Table 7: Numerical blow-up times, numbers of iterations, CPU times (seconds), and orders of the approximations obtained with the explicit Euler method

l	t_n	n	CPU time	s
16	0.090365	7765	15	-
32	0.089372	29778	71	-
64	0.087243	114065	630	0.85
128	0.086658	435971	8580	0.95

Table 8: Numerical blow-up times, numbers of iterations, CPU times (seconds), and orders of the approximations obtained with the implicit Euler method

l	t_n	n	CPU time	s
16	0.090575	7767	19	-
32	0.088423	29780	115	-
64	0.087226	114067	1718	0.85
128	0.086636	435972	23051	1.07

In the following, we also give some plots to illustrate our analysis. In Figures 1 to 8, we can appreciate that the discrete solution blows up in finite time.

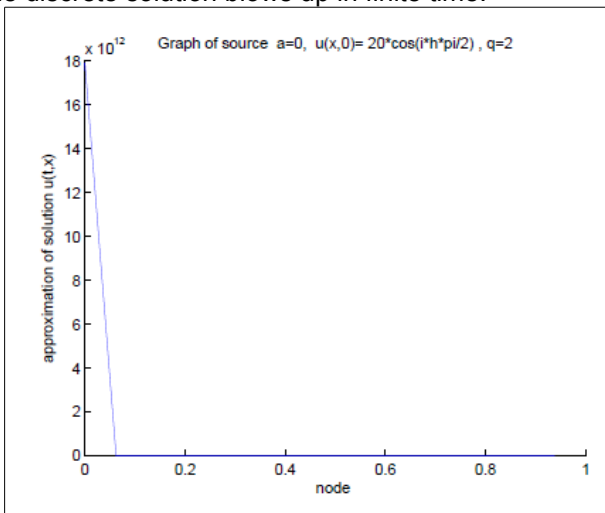


Fig1: Evolution of the discrete solution, source $f(u)=u^q, q=2, a=0$; $l=16$ (implicit scheme).

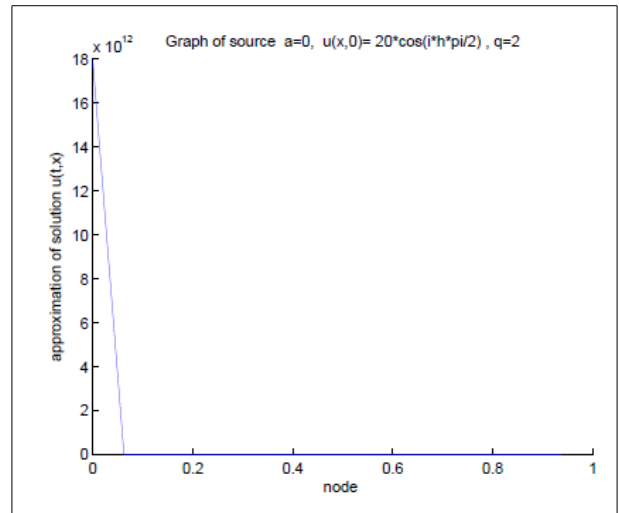


Fig2: Evolution of the discrete solution, source $f(u)=u^q, q=2, a=0$; $l=16$ (explicit scheme).

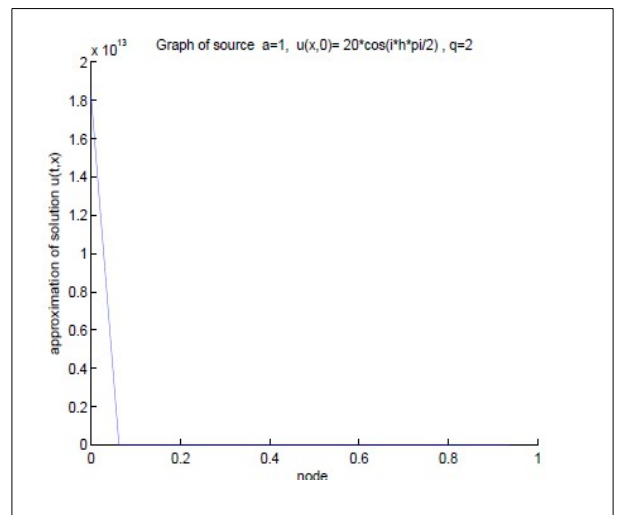


Fig3: Evolution of the discrete solution, source $f(u)=u^q, q=2, a=1$; $l=16$ (implicit scheme).

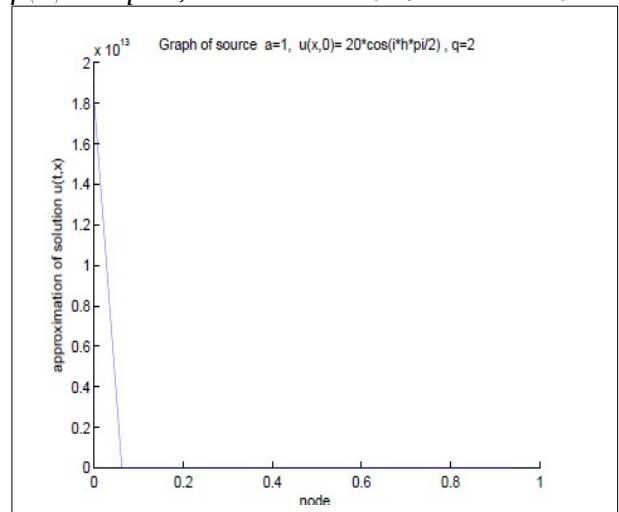


Fig4: Evolution of the discrete solution, source $f(u)=u^q, q=2, a=1$; $l=16$ (explicit scheme).

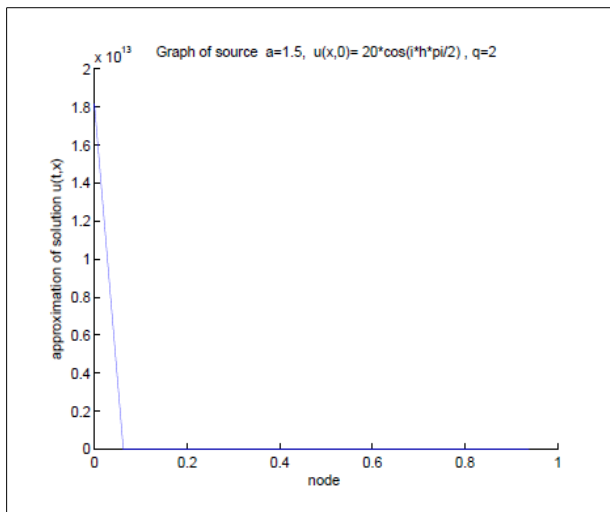


Fig5: Evolution of the discrete solution, source $f(u) = u^q$, $q=2$, $a=1.5$; $l = 16$ (implicit scheme).

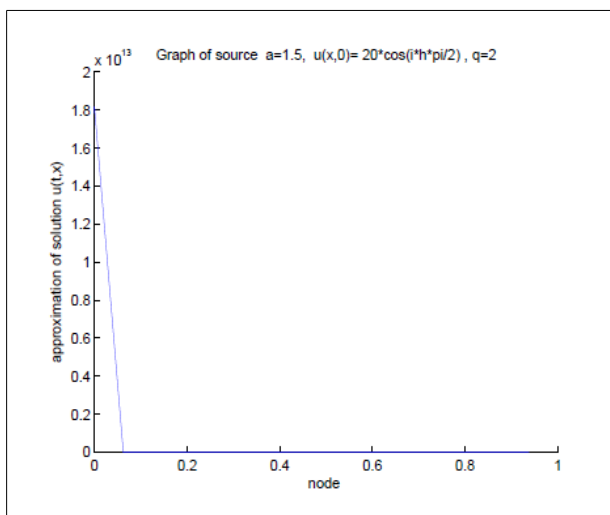


Fig6: Evolution of the discrete solution, source $f(u) = u^q$, $q=2$, $a=1.5$; $l = 16$ (explicit scheme).

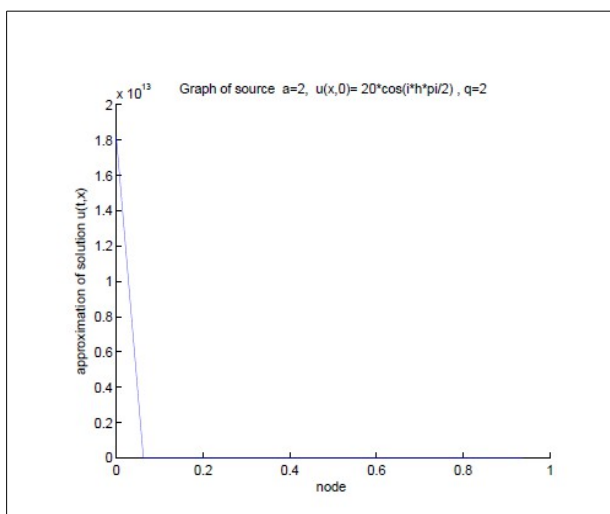


Fig7: Evolution of the discrete solution, source $f(u) = u^q$, $q=2$, $a=2$; $l = 16$ (implicit scheme).

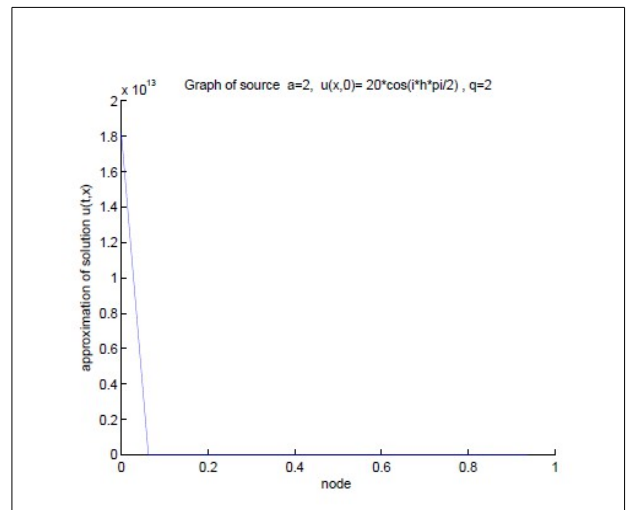


Fig8: Evolution of the discrete solution, source $f(u) = u^q$, $q=2$, $a=2$; $l = 16$ (explicit scheme).

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